

PHYC 3540

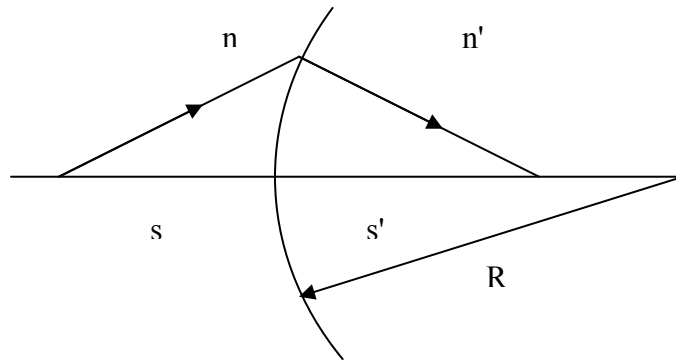
Summary

I. Propagation of light

- a. History, wave vs. particle picture
- b. Fermat's Principle – actual path is that for which the OPL is an extremum
- c. Huygen's principle – every point on wavefront acts as a source
 - i. Large aperture ($\gg \lambda$) – rectilinear propagation (Geometric optics)
 - ii. Small aperture ($\geq \lambda$) – spreading of wavefront (Diffraction)
- d. Wavefronts, Rays – Laws of reflection and refraction (Snell's law)

II. Geometric Optics

- a. Imaging – Cartesian surfaces, Paraxial ray approximation
- b. Refraction at a Spherical interface



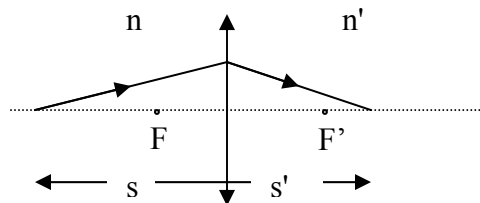
$$\frac{n}{s} + \frac{n'}{s'} = \frac{n' - n}{R} = \frac{n}{f} = \frac{n'}{f'}$$

$$\frac{n}{f} = \frac{n'}{f'} = P \equiv \text{Power}$$

- c. Thin lens

$$\frac{n}{s} + \frac{n'}{s'} = \frac{n_L - n}{R_1} - \frac{n_L - n'}{R_2} = \frac{n}{f} = \frac{n'}{f'} = P$$

Magnification $M = \frac{h'}{h} = -\frac{ns'}{n's}$ for all imaging

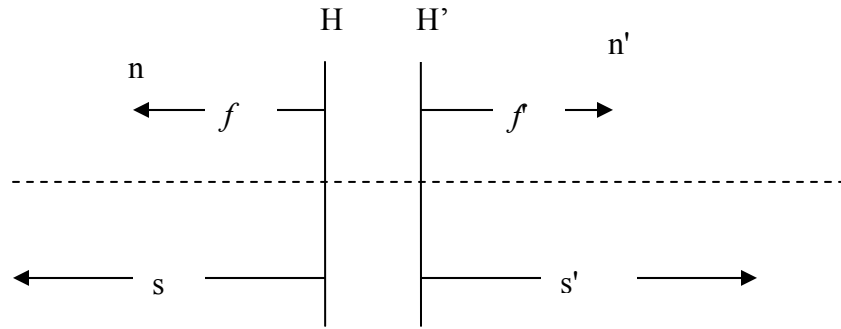


d. Principal Planes (H,H')

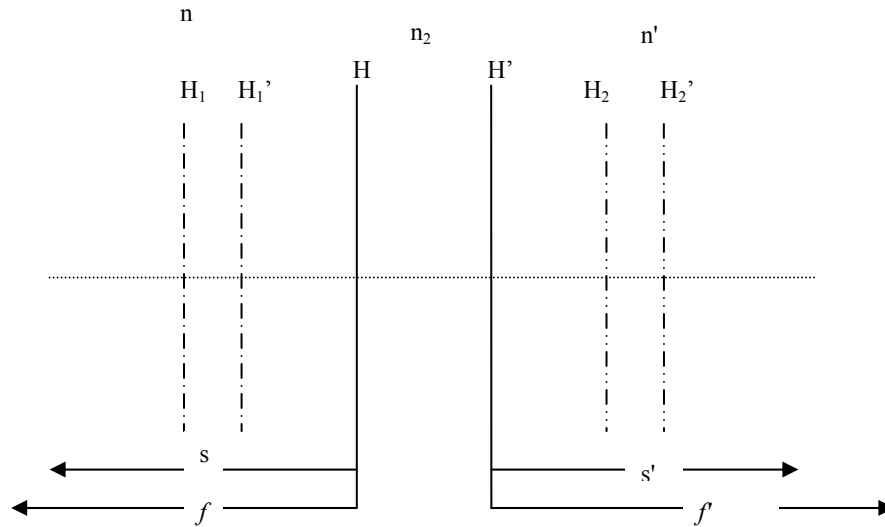
- i. Used to represent any optical element or system and are located such that the system images an object according to:

$$\frac{n}{s} + \frac{n'}{s'} = \frac{n}{f} = \frac{n'}{f'} = P$$

- ii. The principal planes used to represent a spherical interface (or thin lens) must coincide at the interface (or centre of the thin lens).



e. Combination of two systems



- i. Power of combined system $P = P_1 + P_2 - \frac{dP_1P_2}{n_2}$

where, $P_1 = \frac{n}{f_1} = \frac{n_2}{f_1'}$ and $P_2 = \frac{n_2}{f_2} = \frac{n'}{f_2'}$

ii. The location of H, H' are determined with respect to H₁ and H₂' as

$$h' = H_2' H' = -d \left(\frac{P_1}{P} \right) \left(\frac{n'}{n_2} \right) = -d \frac{f'}{f_1'}$$

follows:

$$h = H_1 H = d \left(\frac{P_2}{P} \right) \left(\frac{n}{n_2} \right) = d \frac{f}{f_2}$$

f. Matrix methods in paraxial optics

i. Matrices

$$\text{Translation matrix} \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix}$$

$$\text{Refraction matrix} \begin{bmatrix} 1 & 0 \\ \frac{n-n'}{Rn'} & \frac{n}{n'} \end{bmatrix}$$

$$\text{Reflection matrix} \begin{bmatrix} 1 & 0 \\ \frac{2}{R} & 1 \end{bmatrix}$$

$$\text{Thin lens matrix} \begin{bmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{bmatrix}$$

ii. Cardinal point locations

$$h' = -d \frac{P_1}{P} \left(\frac{n'}{n_2} \right) = \frac{1-A}{C}$$

$$h = d \frac{P_2}{P} \left(\frac{n}{n_2} \right) = \frac{D - (n/n')}{C}$$

g. Stops in optical systems

i. Aperture stop

- Element that limits the bundle of rays (and hence light gathering power or speed) collected by optical system
- Is that element whose entrance pupil subtends the smallest angle at the object point in question
- Exit pupil is the image of the aperture stop in all elements following it
- Entrance pupil is the image of the aperture stop in all elements preceding it

- ii. Field stop
 - Element of the optical system that limits the field of view
 - That element whose entrance window subtends the smallest angle at the centre of the entrance pupil
 - Exit window is the image of the field stop in all elements following it.
 - Entrance window is the image of the field stop in all elements preceding it
- iii. Chief ray, marginal rays – definition of edge of field of view
- iv. Angular field of view
 - Image space – angle formed by the entrance window as viewed from the centre of the entrance pupil
 - Object space – angle formed by the exit window as viewed from the centre of the exit pupil
- h. Applications
 - i. Cameras (aperture settings (f#), exposure, telephoto lens)
 - ii. Human eye – myopia, hyperopia
 - iii. Hand magnifiers, eyepieces, compound microscope
 - iv. Telescope, binoculars

III. **Physical optics** (wave properties are important)

- a. Wave equation – plane waves, spherical waves, cylindrical waves
- b. Amplitude reflection and transmission co-efficients

$$\rho_{12} = \frac{v_2 - v_1}{v_1 + v_2} \quad ; \quad \tau_{12} = \frac{2v_2}{v_1 + v_2}$$

- c. Electromagnetic waves ($E \perp B$; E, B in phase, speed c/n)
 - i. Irradiance

$$I = \frac{1}{2} \epsilon v E_o^2 \quad ; \quad E_o = \text{Amplitude}$$

- ii. Transmittance

$$T = \frac{I_2}{I_1} = 4 \frac{n_1 n_2}{(n_1 + n_2)^2} = \tau_{12} \tau_{21} = \frac{n_2}{n_1} \tau_{12}^2$$

- iii. Reflectance (1→2)

$$R = \frac{I_R}{I_1} = \left(\frac{n_2 - n_1}{n_2 + n_1} \right)^2 = \rho_{12}^2 = \rho_{21}^2$$

- d. Interference

- i. *Two source interference*

- Irradiance

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\phi_2 - \phi_1)$$

$$\text{where, } \phi_2 - \phi_1 = \delta = k(r_2 - r_1) + (\epsilon_2 - \epsilon_1)$$

- If $\epsilon_1 = \epsilon_2$ and $I_1=I_2=I_0$ there are maxima located at,

$$\sin\theta = \pm m \frac{\lambda}{a}$$

- Examples

- Double slit, Lloyds mirror, radio towers
- Dielectric layers, e.g. Haidinger bands, fringes of equal thickness observed in oil films, air wedges
- Antireflection coatings
- Newton's rings

ii. Multiple beam interference (coated dielectric layers)

- Transmitted irradiance

$$I_T = \frac{T^2 I_o}{(1-R)^2} \left(\frac{1}{1 + \frac{4R}{(1-R)^2} \sin^2 \frac{\delta}{2}} \right)$$

where, $\frac{\delta}{2} = k_o n d \cos\theta'$ and n, θ' pertain to the dielectric layer.

- Irradiance maxima occur when
 $2nd \cos\theta' = \pm m \lambda_o$

- Fabry-Perot Interferometer

- Finesse =

$$\mathfrak{F} = \frac{\text{separation between fringes}}{\text{fringe width (FWHM)}} = \frac{\pi \sqrt{R}}{1-R}$$

- Free spectral range

$$\Delta\lambda_{FSR} = \frac{\lambda^2}{2nd} \quad ; \quad \Delta\tilde{\nu}_{FSR} = \frac{1}{2nd}$$

- Fringe width

$$\Delta\lambda = \frac{\text{free spectral range}}{\text{finesse}} = \frac{\lambda^2}{2nd\mathfrak{F}}$$

- Spectral Resolution

$$\mathfrak{R} = \frac{\lambda}{\Delta\lambda_{\min}} = \frac{2nd\mathfrak{F}}{\lambda}$$

$$\mathfrak{R} = \frac{\tilde{\nu}}{\Delta\tilde{\nu}} = 2nd\tilde{\nu}\mathfrak{F}$$

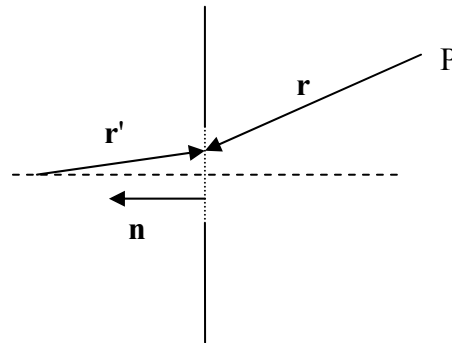
e. Diffraction

- i. Fresnel–Kirchoff diffraction formula (Mathematical statement of Huygen's principle)

$$E = \frac{-ikE_o}{2\pi} \int_{aperture} \frac{e^{ik(r+r')}}{rr'} F(\theta) dA$$

where,

$$F(\theta) = \frac{1}{2} [\cos(\hat{n}, \vec{r}) - \cos(\hat{n}, \vec{r}')] = \textit{obliquity factor}$$



- ii. Fraunhofer limit (plane waves)

diffraction integral reduces to $E = C' \int_{aperture} e^{ikr} dA$

- Single slit (width b)

$$I = I_o \left(\frac{\sin \beta}{\beta} \right)^2 \quad \text{where, } \beta = \frac{kb \sin \theta}{2}$$

$$I = 0 \quad \text{when } \sin \theta = \pm m \frac{\lambda}{b}$$

- Circular aperture (diameter $d_o = 2a_o$)

$$I = I_o \left(\frac{2J_1(\gamma)}{\gamma} \right)^2 \quad \text{where, } \gamma = ka_o \sin \theta$$

$$I = 0 \quad \text{when } \sin \theta = \frac{1.22\lambda}{a_o}$$

Airy disc; Spatial resolution of optical devices; focussing limit

- N slits (width b, separation a)

$$I = I_o \left(\frac{\sin \beta}{\beta} \right)^2 \left(\frac{\sin N\alpha}{\alpha} \right)^2 \quad \text{where, } \alpha = \frac{ka \sin \theta}{2}$$

$$\text{Principal maxima when } \sin \theta = \pm p \frac{\lambda}{a}$$

- Diffraction gratings
Grating dispersion
Spectral resolution $\mathfrak{R} = mN$

f. Polarization

i. Jones vector for polarization

General $\begin{bmatrix} E_{ox} \\ E_{oy}e^{i\delta} \end{bmatrix}$

Vertical $\begin{bmatrix} 0 \\ 1 \end{bmatrix}$ Horizontal $\begin{bmatrix} 1 \\ 0 \end{bmatrix}$

Linearly polarized at an angle $\pm\alpha$, $(-1)^m \begin{bmatrix} \cos\alpha \\ \sin\alpha \end{bmatrix}$

Left Circularly polarized $\frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix}$ (i.e. counterclockwise)

Right Circularly polarized $\frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix}$ (i.e. clockwise)

Left Elliptically polarized $\frac{1}{\sqrt{A^2 + B^2 + C^2}} \begin{bmatrix} A \\ B + iC \end{bmatrix}$

Right Elliptically polarized $\frac{1}{\sqrt{A^2 + B^2 + C^2}} \begin{bmatrix} A \\ B - iC \end{bmatrix}$

where, $E_{ox} = A$ $E_{oy} = \sqrt{B^2 + C^2}$ $\delta = \tan^{-1}\left(\frac{C}{B}\right)$

and the angle of inclination, α , of the ellipse is given by

$$\tan(2\alpha) = \frac{2E_{ox}E_{oy}}{E_{ox}^2 - E_{oy}^2}$$

ii. Jones matrices for optical elements

- Linear polarizer (Transmission axis at an angle θ wrt x-axis)

$$\begin{bmatrix} \cos^2\theta & \sin\theta \cos\theta \\ \sin\theta \cos\theta & \sin^2\theta \end{bmatrix}$$

- Phase retarders

In general, $\begin{bmatrix} e^{i\varepsilon_x} & 0 \\ 0 & e^{i\varepsilon_y} \end{bmatrix}$ where, ε_x and ε_y represent

the advance in phase of the components

- Quarter wave plate ($|\varepsilon_y - \varepsilon_x| = \pi/2$)

$e^{-i\pi/4} \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$ Slow axis vertical

$e^{+i\pi/4} \begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix}$ Slow axis horizontal

➤ Half wave plate ($|\varepsilon_y - \varepsilon_x| = \pi$)

$$e^{-i\pi/2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \text{ Slow axis vertical}$$

$$e^{+i\pi/2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \text{ Slow axis horizontal}$$

- Rotator (light polarized at an angle θ is rotated by an angle β)

$$\begin{bmatrix} \cos \beta & -\sin \beta \\ \sin \beta & \cos \beta \end{bmatrix}$$

g. Holography

i. Hologram

$$R(x, y, z) = r(x, y, z)e^{i\psi(x, y, z)} = re^{ikz}$$

$$O(x, y, z) = o(x, y, z)e^{i\phi(x, y, z)} = oe^{ikr}$$

$$\text{where } r = \sqrt{x^2 + y^2 + z^2}$$

$$I(x, y) \propto |O + R|^2 = OO^* + RR^* + OR^* + O^*R$$

where, R – reference wave, O – object wave, I - intensity

ii. Reconstructed wave

$$R(x, y)t_p = (t_b + BOO^*)R + BR^2O^* + B|R|^2O$$

where, t – transmission, first term – direct wave, second term – conjugate wave, third term – object wave